# FINITE-DIMENSIONAL TOEPLITZ KERNELS AND NEARLY-INVARIANT SUBSPACES 

M.C. CÂMARA and J.R. PARTINGTON

## Communicated by Nikolai K. Nikolski


#### Abstract

A systematic analysis of the structure of finite-dimensional near-ly-invariant subspaces of the Hardy space on the half-plane of index $p$ (with $1<p<\infty$ ) is made, and a criterion given by which they may be recognised. As a consequence, a new approach to Hitt's theorem on nearly-invariant subspaces is developed. Moreover, an analogue is given of Hayashi's theorem for finite-dimensional Toeplitz kernels; this is used to establish a necessary and sufficient condition for a Toeplitz kernel to be non-trivial and of dimension $n$, in terms of a factorisation of its symbol, analogous to Nakazi's work for the disc.


Keywords: Toeplitz operator, Toeplitz kernel, model space, nearly-invariant subspace, inner-outer factorization, Riemann-Hilbert problem.

MSC (2010): 47B35, 30H10.

## 1. INTRODUCTION

Let $H^{p}(\mathbb{D})$ denote the Hardy space on the unit disc $\mathbb{D}$, and let $S^{*}$ denote the backward shift operator defined by

$$
\left(S^{*} f\right)(z)=\frac{f(z)-f(0)}{z}
$$

Nearly $S^{*}$-invariant (abbreviated to n . $S^{*}$-invariant) subspaces of $H^{2}(\mathbb{D})$, i.e., the closed subspaces $\mathcal{E}$ of $H^{2}(\mathbb{D})$ such that $z^{-1} \mathcal{E} \cap H^{2}(\mathbb{D}) \subset \mathcal{E}$, were introduced by Hitt in [17] and have since been the subject of various works ([8], [9], [15], [16], [21], [22]). They are defined analogously in the $H^{p}$ setting, for any $p \in(1, \infty)$, whether we consider the setting of the disc or the upper half-plane. The kernel of a Toeplitz operator $T_{g}$ with an essentially bounded symbol $g$ is a n. $S^{*}$-invariant subspace in each of these settings.

By Hitt's theorem, any n. $S^{*}$-invariant subspace of $H^{2}(\mathbb{D})$ can be described as a product $\mathcal{E}=f K_{\theta}$ where $f$ is the function of unit norm which is orthogonal to
$\mathcal{E} \cap z H^{2}(\mathbb{D})$ in $H^{2}(\mathbb{D})$ and satisfies $f(0)>0$, and $\theta$ is an inner function vanishing at the origin. Here $K_{\theta}$ denotes $H^{2}(\mathbb{D}) \ominus \theta H^{2}(\mathbb{D})$.

In [15] Hayashi addressed the question of which nontrivial closed subspaces of $H^{2}(\mathbb{D})$ are Toeplitz kernels (i.e., kernels of Toeplitz operators) and showed that they were precisely those that could be represented as a product $\mathcal{E}=f K_{\theta}$ where $f \in H^{2}(\mathbb{D})$ with $f^{2}$ rigid in $H^{1}(\mathbb{D})$ and $\theta$ is an inner function vanishing at the origin.

These results were further developed by Sarason in [21] and [22]. Their generalisation to the vectorial case in $H^{2}(\mathbb{D})$ was studied in [9], where a nice introduction to these problems and the corresponding known results can be found.

Few attempts have been made to extend them to the $H^{p}$ setting, possibly due to the lack of a Hilbert space structure allowing one to use a similar line of reasoning. In [13] Dyakonov addressed the question of which closed subspaces of $H^{p}(\mathbb{D}), 1 \leqslant p \leqslant \infty$, are Toeplitz kernels and proposed an alternative parametrisation of the kernel of a Toeplitz operator $T_{g}$ based on Bourgain's factorisation ([1], [3]) for its symbol, trying to cover the whole range $p \in[1, \infty]$ and avoid the use of rigid functions. One disadvantage of this approach is that not only is such a representation highly non-unique, but also it does not indicate the dimension of the Toeplitz kernel and, in particular, whether it is trivial or not.

The question of whether the natural extension of Hitt's and Hayashi's results to the $H^{p}$ setting holds, where naturally $f$ and $\theta$ should depend on $p$, remains open.

In this paper we study the case of finite-dimensional subspaces of $H_{p}^{+}$(using the upper half-plane instead of the disk as in [7]) by taking an approach that may provide some useful lines of reasoning to study the unsolved problem of whether Hitt's and Hayashi's theorems can be extended to all nontrivial subspaces of $H_{p}^{+}$ (or $H^{p}(\mathbb{D})$ ). After some preliminary results and notation, the main results are contained in Sections 2 and 3 . In Section 2 we establish an analogue of Hitt's theorem and we moreover present a simple criterion to recognise a $\mathrm{n} . S^{*}$-invariant subspace in $H_{p}^{+}, p \in(1, \infty)$, by studying the quotients of any two non-zero functions in the subspace. In Section 3 we present an analogue of Hayashi's result for finite-dimensional Toeplitz kernels and use it to establish a necessary and sufficient condition for the kernel of a Toeplitz operator $T_{g}$ to be non-trivial and of dimension $n$, in terms of a factorisation of its symbol. Some of these results provide analogues of the work of Nakazi [20] in the disc, although they are equivalent to Nakazi's results only for $p=2$.

We take $0<p \leqslant \infty$ and $H_{p}^{+}, H_{p}^{-}$to be the Hardy spaces of the upper and lower half-planes $\mathbb{C}^{+}$and $\mathbb{C}^{-}$respectively ( $\left.[12]\right)$. We write $L_{p}$ to denote $L^{p}(\mathbb{R})$. By $\mathcal{G} H_{\infty}^{ \pm}$we denote the class of invertible elements in $H_{\infty}^{ \pm}$, and similarly for $\mathcal{G} L_{\infty}$.

For $G \in L_{\infty}$ and $1<p<\infty$, the Toeplitz operator $T_{G}: H_{p}^{+} \rightarrow H_{p}^{+}$is defined by

$$
T_{G} f_{+}=P^{+}\left(G f_{+}\right), \quad f_{+} \in H_{p}^{+},
$$

where $P^{+}$denotes the projection of $L_{p}$ onto $H_{p}^{+}$parallel to $H_{p}^{-}$. A Toeplitz kernel, or T-kernel, is a subspace of $H_{p}^{+}, 1<p<\infty$, which is the kernel of some Toeplitz operator.
2. NEARLY $S^{*}$-INVARIANT SUBSPACES AND MODEL SPACES

For $1<p<\infty$, and $\theta \in H_{\infty}^{+}$inner, we define the model space $K_{\theta}^{p} \subset H_{p}^{+}$by

$$
K_{\theta}^{p}=H_{p}^{+} \cap \theta H_{p}^{-}
$$

In particular $K_{\theta}^{2}=H_{2}^{+} \ominus \theta H_{2}^{+}$. We shall write $K_{\theta}$ rather than $K_{\theta}^{p}$ when no confusion is likely.

We shall require the functions

$$
\lambda_{ \pm}(\xi)=\xi \pm \mathrm{i} \quad \text { and } \quad r(\xi)=\frac{\xi-\mathrm{i}}{\xi+\mathrm{i}}
$$

and write $S$ for the operator $T_{r}$ on $H_{p}^{+}$of multiplication by $r$, with $S^{*}$ the operator $T_{\bar{r}}$.

DEFINITION 2.1 ([7]). Let $\mathcal{E}$ be a proper closed subspace of $H_{p}^{+}, 1<p<\infty$, and $\eta$ a complex-valued function defined a.e. on $\mathbb{R}$. We say that $\mathcal{E}$ is nearly $\eta$ invariant if, for every $f_{+} \in \mathcal{E}$ such that $\eta f_{+} \in H_{p}^{+}$, we have $\eta f_{+} \in \mathcal{E}$; that is, $\eta \mathcal{E} \cap H_{p}^{+} \subset \mathcal{E}$. If $\mathcal{E}$ is nearly $\eta$-invariant with $\eta \in L_{\infty}$, then we also say that $\mathcal{E}$ is nearly $T_{\eta}$-invariant.

We abbreviate "nearly $\eta$-invariant" to "n. $\eta$-invariant".
Throughout this section $\mathcal{E}$ will denote a n. $S^{*}$-invariant subspace of $H_{p}^{+}$, that is, a nearly $\eta$-invariant subspace for $\eta=\bar{r}$. It is clear that $\mathcal{E}$ is one-dimensional if and only if $\mathcal{E}=\operatorname{span}\left\{f_{+}\right\}$where $f_{+} \in H_{p}^{+}$and $f_{+}(\mathrm{i}) \neq 0$. More generally, we have:

Proposition 2.2. If $\operatorname{dim} \mathcal{E} \geqslant N, N \in \mathbb{N}$, then there are (at least) $N$ linearly independent elements of $\mathcal{E}$ which do not vanish at i .

Proof. If $\mathcal{E}$ is a n. $S^{*}$-invariant subspace of $H_{p}^{+}$, then there exists $f_{1}^{+} \in \mathcal{E}$ such that $f_{1}^{+}(\mathrm{i}) \neq 0$. Let $\left\{f_{1}^{+}, f_{2}^{+}, \ldots, f_{N}^{+}\right\}$be a set of linearly independent elements of $\mathcal{E}$ and define $\widetilde{f}_{j+}:=a_{j} f_{1+}+b_{j} f_{j+}, j=2, \ldots, N$, where $b_{j} \neq 0$ are such that $\widetilde{f}_{j+}(\mathrm{i}) \neq 0$. Then $f_{1+}, \widetilde{f}_{2+}, \ldots, \widetilde{f}_{N+}$ are linearly independent elements of $\mathcal{E}$ that do not vanish at i.

Thus, if $\operatorname{dim} \mathcal{E}=N$, we can choose a basis for $\mathcal{E}$ such that none of its elements vanishes at i.

We also have the following.
Proposition 2.3. If $\operatorname{dim} \mathcal{E} \geqslant N$, with $N \in \mathbb{N}$, then there exists at least one element $\psi_{+} \in \mathcal{E}$ with a zero of order $N-1$ at i .

Proof. In any subspace of $H_{p}^{+}$with $N$ linearly independent elements there exists an element $\psi_{+}$with a zero of order $m \geqslant N-1$ at i. Since $\mathcal{E}$ is n . $S^{*}$-invariant, $r^{-m+N-1} \psi_{+}$belongs to $\mathcal{E}$ and has a zero of order $N-1$ at i.

Remark that if $\operatorname{dim} \mathcal{E}=N$ then the function $\psi_{+}$of Proposition 2.3 is unique up to a constant factor.

As a consequence, we can represent a finite dimensional n. $S^{*}$-invariant subspace in terms of a model space as follows.

THEOREM 2.4. Let $\mathcal{E}$ be a n. $S^{*}$-invariant subspace of $H_{p}^{+}$with $\operatorname{dim} \mathcal{E}=N$ (where $N \in \mathbb{N}$ ) and let $\psi_{+} \in \mathcal{E}$ admit a zero of order $N-1$ at i. Defining $\widetilde{\psi}_{+}:=$ $r^{-N+1} \psi_{+}$, we have

$$
\begin{equation*}
\mathcal{E}=\lambda_{+} \tilde{\psi}_{+} K_{r} N \tag{2.1}
\end{equation*}
$$

Proof. It is clear that $\widetilde{\psi}_{+}, r \widetilde{\psi}_{+}, \ldots, r^{N-1} \tilde{\psi}_{+}$are linearly independent elements of $\mathcal{E}$, so that $\mathcal{E}=\operatorname{span}\left\{\widetilde{\psi}_{+}, r \widetilde{\psi}_{+}, \ldots, r^{N-1} \widetilde{\psi}_{+}\right\}$. Thus $\varphi_{+} \in \mathcal{E}$ if and only if, for some $A_{1}, \ldots, A_{N} \in \mathbb{C}$, we have

$$
\begin{aligned}
\varphi_{+} & =\left(A_{1}+A_{2} r+\cdots+A_{N} r^{N-1}\right) \tilde{\psi}_{+} \\
& =\widetilde{\psi}_{+} \lambda_{+}\left(\frac{A_{1}}{\lambda_{+}}+A_{2} r \frac{1}{\lambda_{+}}+\cdots+A_{N} r^{N-1} \frac{1}{\lambda_{+}}\right)
\end{aligned}
$$

and $\frac{A_{1}}{\lambda_{+}}+A_{2} r \frac{1}{\lambda_{+}}+\cdots+A_{N} r^{N-1} \frac{1}{\lambda_{+}}$is the general form of an element in $K_{r^{N}}$.
The representation (2.1) is unique modulo rational functions belonging to $\mathcal{G} H_{\infty}^{+}$and equivalence of inner functions. To state this more precisely, we introduce here some notation.

DEFINITION 2.5 ([6]). If $\theta_{1}$ and $\theta_{2}$ are inner functions, we say that $\theta_{1} \sim \theta_{2}$ if and only if there are functions $h_{ \pm} \in \mathcal{G} H_{\infty}^{ \pm}$such that

$$
\begin{equation*}
\theta_{1}=h_{-} \theta_{2} h_{+} \tag{2.2}
\end{equation*}
$$

and we say that $K_{\theta_{1}} \sim K_{\theta_{2}}$ if and only if

$$
\begin{equation*}
K_{\theta_{1}}=h_{+} K_{\theta_{2}} \quad \text { with } h_{+} \in \mathcal{G} H_{\infty}^{+} \tag{2.3}
\end{equation*}
$$

It is clear that $\theta_{1} \sim \theta_{2} \Rightarrow K_{\theta_{1}} \sim K_{\theta_{2}}$ ([6]). We will use the notation

$$
\begin{equation*}
\frac{K_{\theta_{1}}}{K_{\theta_{2}}} \simeq h_{+} \tag{2.4}
\end{equation*}
$$

whenever (2.3) holds.
THEOREM 2.6. Let $\mathcal{E} \subset H_{p}^{+}$be a $n$. $S^{*}$-invariant subspace with dimension $N \in$ $\mathbb{N}$. If, for some function $F_{+} \in H_{p}^{+}$and an inner function $\theta$, we have

$$
\begin{equation*}
\mathcal{E}=\lambda_{+} F_{+} K_{\theta} \tag{2.5}
\end{equation*}
$$

then $F_{+}(\mathrm{i}) \neq 0, \theta \sim r^{N}$ and, for any function $\tilde{\psi}_{+}$satisfying (2.1),

$$
\begin{equation*}
\frac{K_{\theta}}{K_{r^{N}}} \simeq \frac{\widetilde{\psi}_{+}}{F_{+}} \in \mathcal{R} \cap \mathcal{G} H_{\infty}^{+} \tag{2.6}
\end{equation*}
$$

where $\mathcal{R}$ denotes the set of all rational functions in $L_{\infty}$.
Proof. It is clear that $F^{+}(\mathrm{i}) \neq 0$, by Proposition 2.2. On the other hand, since $\operatorname{dim} K_{\theta}=\operatorname{dim} \mathcal{E}=N, \theta$ is a rational inner function and therefore $\theta=h_{-} r^{N} h_{+}$ where $h_{ \pm} \in \mathcal{G} H_{\infty}^{ \pm}$are rational functions. It follows that $K_{\theta}=h_{+} K_{r N}$ and from (2.1) and (2.5) we have

$$
\begin{equation*}
\mathcal{E}=\lambda_{+} \widetilde{\psi}_{+} K_{r N}=\lambda_{+} F_{+} h_{+} K_{r N} . \tag{2.7}
\end{equation*}
$$

Since the second equality in (2.7) implies that $F_{+} h_{+} \in \mathcal{E}$, it follows from the first equality in 2.7) that, for some constants $A_{1}, A_{2}, \ldots, A_{N} \in \mathbb{C}$, with $A_{1} \neq 0$,

$$
\begin{equation*}
F_{+} h_{+}=\widetilde{\psi}_{+}\left(A_{1}+A_{2} r+\cdots+A_{N} r^{N-1}\right) \tag{2.8}
\end{equation*}
$$

For $N=1$ this means that $F_{+} h_{+}=A_{1} \tilde{\psi}_{+}$with $A_{1} \neq 0$ and therefore

$$
\begin{equation*}
\frac{\tilde{\psi}_{+}}{F_{+}}=\frac{h_{+}}{A_{1}}, \tag{2.9}
\end{equation*}
$$

so that 2.6 holds. If $N>1$ it follows from (2.7) that we also have

$$
\begin{equation*}
r^{j+1} F_{+} h_{+} \in \mathcal{E} \quad \text { for all } j=0,1, \ldots, N-2 \tag{2.10}
\end{equation*}
$$

and, taking 2.8 into account,

$$
\begin{equation*}
r^{j+1} F_{+} h_{+}=\widetilde{\psi}_{+}\left(A_{1} r^{j+1}+\cdots+A_{N-j-1} r^{N-1}\right)+\widetilde{\psi}_{+}\left(A_{N-j} r^{N}+\cdots+A_{N} r^{N+j}\right) \tag{2.11}
\end{equation*}
$$

Since, by (2.10), the left hand side of this equality represents a function in $\mathcal{E}$ and $\widetilde{\psi}_{+}\left(A_{1} r^{j+1}+\cdots+A_{N-j-1} r^{N-1}\right) \in \lambda_{+} \widetilde{\psi}_{+} K_{r N}=\mathcal{E}$, we see that 2.11) implies that

$$
\eta_{j}:=\widetilde{\psi}_{+}\left(A_{N-j} r^{N}+\cdots+A_{N} r^{N+j}\right) \in \mathcal{E}
$$

and therefore, since $\eta_{j}$ has a zero of order greater or equal to $N$ at i, $\eta_{j}=0$ for all $j=0,1, \ldots, N-2$. We have thus

$$
A_{N-j} r^{N}+\cdots+A_{N} r^{N+j}=0 \quad \text { for all } j=0,1, \ldots, N-2
$$

and it follows that $A_{N}=A_{N-1}=\cdots=A_{2}=0$. From 2.8 we see therefore that 2.9 and, consequently, (2.6) hold.

Defining $\varphi_{+}^{*}=\frac{\widetilde{\psi}_{+}}{\lambda_{+}^{N-1}}$, and noting that the set of functions

$$
\left\{\lambda_{+}^{N-1}, \lambda_{+}^{N-2} \lambda_{-}, \ldots, \lambda_{+} \lambda_{-}^{N-2}, \lambda_{-}^{N-1}\right\}
$$

forms a basis for the space $\mathcal{P}_{N-1}$ of all polynomials of degree at most $N-1$, we arrive at the following result.

THEOREM 2.7. Let $\mathcal{E}$ be a $n . S^{*}$-invariant subspace of $H_{p}^{+}$with $\operatorname{dim} \mathcal{E}=N$. Then there is a function $\varphi_{+}^{*} \in \mathcal{E} \backslash r H_{p}^{+}$such that

$$
\begin{equation*}
\mathcal{E}=\left\{\varphi_{+}^{*} p_{+}: p_{+} \in \mathcal{P}_{N-1}\right\} . \tag{2.12}
\end{equation*}
$$

In the case that $\mathcal{E}$ is a Toeplitz kernel, and hence nearly-invariant under division by every inner function ([7]), we may also conclude that $\varphi_{+}^{*}$ is outer. In fact, a nontrivial Toeplitz kernel cannot be contained in $\theta H_{p}^{+}$if $\theta$ is a non constant inner function ([6], Theorem 2.4).

We may ask then if, conversely, any set of the form (2.12) is a n. $S^{*}$-invariant subspace of $H_{p}^{+}$and, in case $\varphi_{+}^{*}$ is outer, if it is a Toeplitz kernel. While the latter question will be dealt with in the next section, the answer to the former is given in the following theorem, which moreover provides a simple criterion to recognise a finite-dimensional $\mathrm{n} . S^{*}$-invariant subspace of $H_{p}^{+}, p \in(1, \infty)$.

THEOREM 2.8. Suppose that $\mathcal{E} \subset H_{p}^{+}$and $\operatorname{dim} E=N$. Then $\mathcal{E}$ is n. $S^{*}$-invariant if and only if:
(i) $\mathcal{E}$ contains at least one function that does not vanish at $i$, and
(ii) the quotient of any two functions in $\mathcal{E}$ is equal to a quotient of two polynomials of degrees at most $N-1$.

Proof. The conditions are obviously necessary, by Theorem 2.7 .
To show their sufficiency, we may clearly take $N>1$. Pick a basis $v_{1}^{+}, \ldots, v_{N}^{+}$ of $\mathcal{E}$, and assume that $v_{1}^{+}(\mathrm{i}) \neq 0$, as we can by Proposition 2.2

For each $k=1, \ldots, N$ we write

$$
v_{k}^{+}=\frac{p_{k}}{q_{k}} v_{1}^{+}
$$

where $p_{1}=q_{1}=1$, and in general $p_{k}, q_{k}$ are polynomials of degree at most $N-1$ with no common factors. Since $v_{1}^{+}(\mathrm{i}) \neq 0$, we have $q_{k}(\mathrm{i}) \neq 0$ for all $k=2, \ldots, N$.

Let $d$ denote the least common multiple of the polynomials $q_{k}, k=2, \ldots, N$. Clearly the set of zeroes of $d$ is finite and does not contain i. We claim that there exists a complex linear combination

$$
r=\frac{c_{2} v_{2}^{+}+\cdots+c_{N} v_{N}^{+}}{v_{1}^{+}}
$$

such that every zero of multiplicity $m$ in $d$ is a pole of order $m$ in $r$, which implies that $d$ has at most $N-1$ zeroes, counting multiplicity.

For suppose that $z_{0}$ is a zero of $d$ with multiplicity $m$. Then, for some $k_{0} \in$ $\{2, \ldots, N\}, z_{0}$ must be a zero of $q_{k_{0}}$ with multiplicity $m$ and, in a neighbourhood of $z_{0}$, we have

$$
\frac{v_{k}^{+}(\xi)}{v_{1}^{+}(\xi)}=\frac{1}{\left(\xi-z_{0}\right)^{m}}\left(b_{k}+O\left(\xi-z_{0}\right)\right), \quad k=2, \ldots, N,
$$

where at least one $b_{k}$ is non-zero. The point $z_{0}$ will be a pole of order $m$ in $r$ if and only if $\sum_{k=2}^{N} c_{k} b_{k} \neq 0$. Repeating the same reasoning for all zeroes of $d$, we see that it is sufficient to choose a point $\left(c_{2}, \ldots, c_{N}\right)$ in $\mathbb{C}^{N-1}$ which does not belong to finitely-many hyperplanes of the form $\sum_{k=2}^{N} c_{k} b_{k}=0$.

Our conclusion is that all the $\frac{v_{k}^{+}}{v_{1}^{+}}$can be written over a common denominator $d \in \mathcal{P}_{N-1}$. Thus

$$
\mathcal{E}=\left\{v_{1}^{+} \sum_{k=1}^{N} \frac{\alpha_{k} \widetilde{p}_{k}}{d}: \alpha_{1}, \ldots, \alpha_{N} \in \mathbb{C}\right\}
$$

for some polynomials $\widetilde{p}_{1}, \ldots, \widetilde{p}_{N}$ belonging to $\mathcal{P}_{N-1}$, and indeed we may take $\tilde{p}_{1}=d$. Since $\operatorname{dim} \mathcal{E}=N$, we see that

$$
\mathcal{E}=\left\{v_{1}^{+} \frac{Q}{d}: Q \in \mathcal{P}_{N-1}\right\}
$$

Now near-invariance is clear: since $\mathrm{d}(\mathrm{i}) \neq 0$, if a function $f \in \mathcal{E}$ vanishes at i , then so does the corresponding polynomial $Q$, and thus $\frac{f}{r} \in \mathcal{E}$.

The results of Theorems 2.4 and 2.6 naturally lead to the question of how they relate to Hitt's characterisation of n. $S^{*}$-invariant subspaces [17], [21] in the case when $p=2$ and $n \neq 1$, since in general $\widetilde{\psi}_{+}$is not orthogonal to the space $\mathcal{E}_{0}$ given by

$$
\begin{equation*}
\mathcal{E}_{0}:=\mathcal{E} \cap r H_{2}^{+}=\operatorname{span}\left\{\psi_{+}, r^{-1} \psi_{+}, \ldots, r^{-(N-2)} \psi_{+}\right\}, \tag{2.13}
\end{equation*}
$$

using the notation of Theorem 2.4 .
We can nevertheless obtain an element $\varphi_{+} \in \mathcal{E} \cap \mathcal{E}_{0}^{\perp}$ by Gram-Schmidt orthogonalisation of the basis

$$
\left\{\psi_{+}, r^{-1} \psi_{+}, \ldots, r^{-(N-2)} \psi_{+}, r^{-(N-1)} \psi_{+}\right\}
$$

where $r^{-(N-1)} \psi_{+}=\tilde{\psi}_{+}$as in Theorem 2.4. which yields an orthogonal basis $\left\{\varphi_{1+}, \ldots, \varphi_{N+}\right\}$.

In particular we have

$$
\begin{equation*}
\psi_{+}=\varphi_{1+} \quad \text { and } \quad \varphi_{+}:=\varphi_{N+} \in \mathcal{E} \ominus \mathcal{E}_{0} \tag{2.14}
\end{equation*}
$$

Since $\mathcal{E}_{0}$ has codimension 1 in $\mathcal{E}$, the orthogonal element $\varphi_{+}$is unique apart from a constant factor. From (2.1) it follows that, for some constants $A_{0}, A_{1}, \ldots, A_{N-1} \in$ $\mathbb{C}$, with $A:=A_{N-1} \neq 0$, and $C_{1}, C_{2}, \ldots, C_{N-1} \in \mathbb{C}$ we have

$$
\begin{aligned}
\varphi_{+} & =A_{0} \psi_{+}+A_{1} r^{-1} \psi_{+}+\cdots+A_{N-1} r^{-(N-1)} \psi_{+} \\
& =A\left(r^{-1}-C_{1}\right)\left(r^{-1}-C_{2}\right) \cdots\left(r^{-1}-C_{N-1}\right) \psi_{+} \\
& =A\left(r^{-1}-C_{j}\right) \prod_{s=1, s \neq j}^{N-1}\left(r^{-1}-C_{s}\right) \psi_{+}
\end{aligned}
$$

for any $j=1,2, \ldots, N-1$. So, defining

$$
\psi_{j+}^{*}=\prod_{s=1, s \neq j}^{N-1}\left(r^{-1}-C_{s}\right) \psi_{+}
$$

we then have

$$
\begin{equation*}
\varphi_{+}=A r^{-1} \psi_{j+}^{*}-A C_{j} \psi_{j+}^{*} \tag{2.15}
\end{equation*}
$$

Now $\varphi_{+} \in \mathcal{E}_{0}^{\perp}$ and $\psi_{j+}^{*} \in \mathcal{E}_{0}$, so $\left\langle\varphi_{+}, \psi_{j+}^{*}\right\rangle=0$ and thus the constant $C_{j}$ in 2.15 is given by

$$
C_{j}=\frac{\left\langle r^{-1} \psi_{j+}^{*}, \psi_{j+}^{*}\right\rangle}{\left\langle\psi_{j+}^{*}, \psi_{j+}^{*}\right\rangle}
$$

Since $\left\|r^{-1} \psi_{j+}^{*}\right\|_{2}=\left\|\psi_{j+}^{*}\right\|_{2}$, it follows from the fact that the Cauchy-Schwarz inequality is strict unless the vectors involved are linearly dependent that $\left|C_{j}\right|<1$.

Now we may write

$$
\begin{equation*}
r^{-1}(\xi)-C_{j}=\frac{\xi+\mathrm{i}}{\xi-\mathrm{i}}-C_{j}=\frac{\left(1-C_{j}\right)\left(\xi-\xi_{j}\right)}{\xi-\mathrm{i}} \tag{2.16}
\end{equation*}
$$

where $\xi_{j}=-i \frac{1+C_{j}}{1-C_{j}} \in \mathbb{C}^{-}$since $\left|C_{j}\right|<1$. We have thus proved the following.
ThEOREM 2.9. Let $\varphi_{+} \in \mathcal{E} \ominus \mathcal{E}_{0}$ with $\varphi_{+} \neq 0$. Then

$$
\begin{equation*}
\varphi_{+}(\xi)=A \frac{\left(\xi-\xi_{1}\right) \cdots\left(\xi-\xi_{N-1}\right)}{(\xi-\mathrm{i})^{N-1}} \psi_{+}(\xi) \tag{2.17}
\end{equation*}
$$

where $A \in \mathbb{C} \backslash\{0\}, \xi_{1}, \ldots, \xi_{N-1} \in \mathbb{C}^{-}$and $\psi_{+} \in \mathcal{E}$ has a zero of order $N-1$ at i .
It follows from Theorems 2.4 and 2.9 that $\mathcal{E}=\lambda_{+} \tilde{\psi}_{+} K_{r^{N}}=\lambda_{+} \varphi_{+} h_{+} K_{r N}$, where $h_{+} \in \mathcal{G} H_{\infty}^{+}$is given by

$$
h_{+}(\xi)=\frac{(\xi+\mathrm{i})^{N-1}}{\left(\xi-\xi_{1}\right) \cdots\left(\xi-\xi_{N-1}\right)}
$$

Now

$$
h_{+} K_{r^{N}}=h_{+} \operatorname{ker} T_{r^{-N}}=\operatorname{ker} T_{h_{+}^{-1} r^{-N}}=\operatorname{ker} T_{g_{-} \bar{B}^{\prime}},
$$

where $g_{-} \in \mathcal{G} H_{\infty}^{-}$is given by

$$
g_{-}(\xi)=\frac{\left(\xi-\bar{\xi}_{1}\right) \cdots\left(\xi-\bar{\xi}_{N-1}\right)}{(\xi-i)^{N-1}}
$$

and $B$ is the Blaschke product given by

$$
\begin{equation*}
B=\frac{\xi-\mathrm{i}}{\xi+\mathrm{i}} \frac{\xi-\bar{\xi}_{1}}{\bar{\xi}-\bar{\xi}_{1}} \cdots \frac{\xi-\bar{\xi}_{N-1}}{\xi-\bar{\xi}_{N-1}} . \tag{2.18}
\end{equation*}
$$

Since $\operatorname{ker} T_{g_{-} \bar{B}}=\operatorname{ker} T_{\bar{B}}=K_{B}$, we have established the following theorem.
 $\psi_{+}$be the (unique, up to a constant factor) element of $\mathcal{E}$ admitting a zero of order $N-1$ at i , and let $\varphi_{+} \in \mathcal{E} \ominus \mathcal{E}_{0}$. Then

$$
\mathcal{E}=\lambda_{+} \varphi_{+} K_{B},
$$

where $B$ is the finite Blaschke product given by (2.18, where $\xi_{1}, \ldots, \xi_{N-1}$ are the zeroes of the rational function $\frac{\varphi_{+}}{\psi_{+}}$.

Besides establishing a clear relation between Theorem 2.4 and Hitt's theorem, this result moreover defines explicitly the model space associated with Hitt's representation.

## 3. ON FINITE-DIMENSIONAL TOEPLITZ KERNELS

Next we address two closely related questions: when does a Toeplitz operator have a nontrivial kernel of finite dimension, and when is a finite-dimensional subspace of $H_{p}^{+}$a T-kernel?

Here we need the theory of rigid functions.
DEFINITION 3.1. A function $f_{+} \in H_{q}^{+} \backslash\{0\}$, with $0<q<\infty$, is called rigid if and only if, for any $g_{+} \in H_{q}^{+}$such that $\frac{g_{+}}{f_{+}}>0$ a.e. on $\mathbb{R}$, we have $g_{+}=\lambda f_{+}$for some $\lambda \in \mathbb{R}^{+}$.

A rigid function is outer (in $H_{q}^{+}$), and every rigid function in $H_{q}^{+}$is the square of an outer function in $H_{p}^{+}$, with $p=2 q$ ([7], [22]). If $f_{+} \in H_{p}^{+}$and $f_{+}^{2}$ is rigid in $H_{p / 2}^{+}$, we say that $f_{+}$is square-rigid in $H_{p}^{+}$.

It was shown in [22] that a one-dimensional subspace of $H^{2}(\mathbb{D})$ is a $T$-kernel if and only if it is spanned by a function that is square-rigid in $H^{2}(\mathbb{D})$. An analogous result holds for one-dimensional subspaces of $H_{p}^{+}$, for $1<p<\infty$, as follows.

THEOREM 3.2 ([7]). Let $f_{+} \in H_{p}^{+}, 1<p<\infty$. Then span $\left\{f_{+}\right\}$is a T-kernel in $H_{p}^{+}$if and only if $f_{+}$is outer and square-rigid in $H_{p}^{+}$. In that case span $\left\{f_{+}\right\}=$ $\operatorname{ker} T_{\bar{f}_{+} / f_{+}}$.

As a consequence, we also have:
COROLLARy 3.3. If $O_{+} \in H_{p}^{+}$is outer and square-rigid then, for every $k \in \mathbb{N}$,

$$
\begin{equation*}
\operatorname{ker} T_{r^{-k} \bar{O}_{+} / O_{+}}=\operatorname{span}\left\{O_{+}, r O_{+}, \ldots, r^{k} O_{+}\right\}=\lambda_{+} O_{+} K_{r^{k+1}} \tag{3.1}
\end{equation*}
$$

Proof. It is clear that $O_{+}, r O_{+}, \ldots, r^{k} O_{+}$are linearly independent and belong to ker $T_{r^{-k}} \bar{O}_{+} / O_{+}$. If the dimension of the latter was greater than $k+1$ we would have ([2])

$$
\operatorname{ker} T_{\bar{O}_{+} / O_{+}}=\operatorname{ker} T_{r^{k}(r-k}{\left.\overline{O_{+}} / O_{+}\right)}>1,
$$

which is false since $\operatorname{ker} T_{\bar{O}_{+} / O_{+}}=\operatorname{span}\left\{O_{+}\right\}$by Theorem 3.2 .
The following theorem is an analogue of Hayashi's result ([15]) for finitedimensional subspaces of $H_{p}^{+}$.

THEOREM 3.4. Suppose that $\mathcal{E} \subset H_{p}^{+}$and $\operatorname{dim} E=N$. Then $\mathcal{E}$ is a $T$-kernel if and only if $\mathcal{E}=\lambda_{+} O_{+} K_{r^{N}}$, with $O_{+}$a square-rigid outer function in $H_{p}^{+}$.

Proof. Let $\mathcal{E}=\lambda_{+} \varphi_{+} K_{r^{N}}$, according to Theorem 2.4. If $\mathcal{E}$ is a $T$-kernel, then $\varphi_{+}$must be outer, as remarked in Section 2 , so we may write $\varphi_{+}=O_{+}$, with $O_{+}$ outer. Since $\mathcal{E}$ is a $T$-kernel containing $O_{+}$, it must contain the minimal kernel $\mathcal{K}_{\text {min }}\left(O_{+}\right)=\operatorname{ker} T_{\bar{O}_{+} / O_{+}}$.

Let $\psi_{+} \in \operatorname{ker} T_{\bar{O}_{+} / O_{+}} \subset \mathcal{E}$; then $\psi_{+}=\lambda_{+} O_{+} k_{+}$, where $k_{+} \in K_{r N}$ and, for some $\psi_{-} \in H_{p}^{-}$, we have $\frac{\bar{\sigma}_{+}}{O_{+}} \psi_{+}=\psi_{-}$. On the other hand,

$$
\frac{\bar{O}_{+}}{O_{+}} \psi_{+}=\psi_{-} \Longleftrightarrow \bar{O}_{+} \lambda_{+} k_{+}=\psi_{-} \Longleftrightarrow \frac{\lambda_{+}}{\lambda_{-}} k_{+}=\frac{\lambda_{-}^{-1} \psi_{-}}{\bar{O}_{+}}
$$

and, since $\frac{\lambda_{-}^{-1} \psi_{-}}{\bar{O}_{+}} \in L_{p} \cap \overline{\mathcal{N}}_{+}$, we have $\frac{\lambda_{-}^{-1} \psi_{-}}{\bar{O}_{+}} \in H_{p}^{-}$, i.e., $\frac{\psi_{-}}{\bar{O}_{+}} \in \lambda_{-} H_{p}^{-}$.
From

$$
\underbrace{\lambda_{+} k_{+}}_{\in \lambda_{+} H_{p}^{+}}=\underbrace{\frac{\psi_{-}}{\bar{O}_{+}}}_{\in \lambda_{-} H_{p}^{-}}
$$

it follows that both sides are constant ([23]), so that $\psi_{-}=c \bar{O}_{+}$and $\psi_{+}=c O_{+}$ with $c \in \mathbb{C}$. Thus ker $T_{\bar{O}_{+} / O_{+}}=\operatorname{span}\left\{O_{+}\right\}$, which is the same as saying that $O_{+}^{2}$ is rigid in $H_{p / 2}^{+}$.

Conversely, assume that $\mathcal{E}=O_{+} \lambda_{+} K_{r_{N}}$ with $O_{+}^{2}$ rigid in $H_{p / 2}^{+}$. Then, by Corollary 3.3 .

$$
\mathcal{E}=\lambda_{+} O_{+} K_{r^{N}}=\operatorname{ker} T_{r^{-(N-1)}\left(\bar{O}_{+} / O_{+}\right)^{\prime}}
$$

so $\mathcal{E}$ is a $T$-kernel.
Using the result of Theorem 3.4 we can also characterise any non zero finitedimensional $T$-kernel and establish conditions for a $T$-kernel to be trivial, in terms of a factorisation of the symbol of the corresponding Toeplitz operator.

We will need the following results.
THEOREM 3.5 (|7|). For every $\varphi_{+} \in H_{p}^{+} \backslash\{0\}$ there exists a $T$-kernel containing $\varphi_{+}$, denoted by $\mathcal{K}_{\min }\left(\varphi_{+}\right)$, such that for any $g \in L_{\infty}$ we have

$$
\begin{equation*}
\varphi_{+} \in \operatorname{ker} T_{g} \Rightarrow \mathcal{K}_{\min }\left(\varphi_{+}\right) \subset \operatorname{ker} T_{g} \tag{3.2}
\end{equation*}
$$

and, if $\varphi_{+}=I_{+} O_{+}$is an inner-outer factorisation of $\varphi_{+}$,

$$
\begin{equation*}
\mathcal{K}_{\min }\left(\varphi_{+}\right)=\operatorname{ker} T_{\bar{I}_{+} \bar{O}_{+} / O_{+}} \tag{3.3}
\end{equation*}
$$

$\mathcal{K}_{\min }\left(\varphi_{+}\right)$is called the minimal kernel for $\varphi_{+}$. It can be shown moreover that a nontrivial, proper, n . $S^{*}$-invariant subspace $\mathcal{E}$ of $H_{p}^{+}(1<p<\infty)$ is a $T$-kernel if and only if there exists $\varphi_{+} \in H_{p}^{+}$such that $\mathcal{E}=\mathcal{K}_{\min }\left(\varphi_{+}\right)$([7]).

DEFINITION $3.6\left([\sqrt{6})\right.$. If $K=\mathcal{K}_{\min }\left(\varphi_{+}\right)$, we say that $\varphi_{+}$is a maximal function for $K$.

Clearly, if $\varphi_{+}$is a maximal function for $\operatorname{ker} T_{g}$, then we have $g \varphi_{+}=\varphi_{-}$ where $\varphi_{-} \in H_{p}^{-}$is outer ([6]).

THEOREM 3.7. For $g \in L_{\infty}, \operatorname{ker} T_{g}$ is nontrivial and of finite dimension if and only if, for some $N \in \mathbb{N}$, g admits a factorisation

$$
\begin{equation*}
g=g_{-} r^{-N} g_{+}^{-1} \tag{3.4}
\end{equation*}
$$

where $\frac{g_{-}}{\lambda_{-}} \in H_{p}^{-}$is outer and $\frac{g_{+}}{\lambda_{+}} \in H_{p}^{+}$is outer and square-rigid. In that case $\operatorname{ker} T_{g}=$ $\operatorname{ker} T_{r^{-N}\left(\bar{g}_{+} / g_{+}\right)}$and $\operatorname{dim} \operatorname{ker} T_{g}=N$.

Proof. Step 1. Let $O_{+}=\frac{g_{+}}{\lambda_{+}}, O_{-}=\frac{g_{-}}{\lambda_{-}}$. If $g$ admits a representation of the form (3.4) then $\left\{O_{+}, r O_{+}, \ldots, r^{N-1} O_{+}\right\} \subset \operatorname{ker} T_{g}$ and so, by Corollary 3.3. $\operatorname{ker} T_{g} \supset \operatorname{ker} T_{r^{-N+1}\left(\bar{O}_{+} / O_{+}\right)} \neq\{0\}$. On the other hand, if $\varphi_{+} \in \operatorname{ker} T_{g}$, we have $g \varphi_{+}=\varphi_{-}$with $\varphi_{-} \in H_{p}^{-}$which, taking (3.4) into account, is equivalent to

$$
\begin{equation*}
r^{-N+1} \frac{\bar{O}_{+}}{O_{+}} \varphi_{+}=\frac{\bar{O}_{+} \varphi_{-}}{O_{-}} \tag{3.5}
\end{equation*}
$$

Since the right hand side of (3.5) represents a function whose conjugate belongs to the Smirnov class $\mathcal{N}_{+}$as well as to $L_{p}$, it is a function in $H_{p}^{-}$; it follows that $\varphi_{+} \in \operatorname{ker} T_{r^{-N+1}\left(\bar{O}_{+} / O_{+}\right)}$. Therefore (3.4) implies that $\operatorname{ker} T_{g} \subset \operatorname{ker} T_{r^{-N+1}\left(\bar{O}_{+} / O_{+}\right)}$. It follows that $\operatorname{ker} T_{g}=\operatorname{ker} T_{r^{-N+1}\left(\bar{O}_{+} / O_{+}\right)}=\lambda_{+} O_{+} K_{r^{N}}$ and dim $\operatorname{ker} T_{g}=N$.

Step 2. Conversely, let now ker $T_{g}$ be a (n. $S^{*}$-invariant) subspace of $H_{p}^{+}$ with dimension $N$. Let $\varphi_{+}$be a maximal function for $\operatorname{ker} T_{g}$. By Theorem 3.4. $\operatorname{ker} T_{g}=O_{+} \lambda_{+} K_{r^{N}}$, where $O_{+} \in \operatorname{ker} T_{g}$ is outer and square-rigid, so we have

$$
\begin{equation*}
r^{-N} O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+}=\psi_{-} \tag{3.6}
\end{equation*}
$$

where $\psi_{-} \in H_{p}^{-}$is outer by Lemma 3.8 below. On the other hand,

$$
\begin{equation*}
g \varphi_{+}=O_{-} \tag{3.7}
\end{equation*}
$$

where $O_{-}$is outer since $\varphi_{+}$is maximal in ker $T_{g}$. From (3.6) and 3.7) we obtain

$$
g=\frac{O_{-}}{\psi_{-}} r^{-N} \frac{O_{+}^{-1}}{\lambda_{+}}
$$

where $\frac{O_{-}}{\psi_{-} \lambda_{-}}=g r^{-N+1} O_{+} \in H_{p}^{-}$is outer and $\left(\frac{O_{+}^{-1}}{\lambda_{+}}\right)^{-1} \lambda_{+}^{-1}=O_{+}$is outer and square-rigid.

LEMMA 3.8. For $g \in L_{\infty}$ let $\varphi_{+}$be a maximal function in $\operatorname{ker} T_{g}=O_{+} \lambda_{+} K_{r} N$, $N \in \mathbb{N}$, as in Theorem 3.4 Then $\psi_{-}=r^{-N} O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+}$is outer in $H_{p}^{-}$.

Proof. Since $\varphi_{+} \in \operatorname{ker} T_{g}$, we have $\varphi_{+}=O_{+} \lambda_{+} h_{+}$with $h_{+} \in \operatorname{ker} T_{r^{-N}}=$ $K_{r^{N}}$ and thus

$$
\begin{equation*}
r^{-N} O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+}=\psi-\in H_{p}^{-} \tag{3.8}
\end{equation*}
$$

Remark that $O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+} \in \mathcal{N}_{+} \cap L_{p}$, so that $O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+} \in H_{p}^{+}$. If $\psi_{-}=I_{-} O_{-}$is an inner-outer factorisation of $\psi_{-}$in $H_{p}^{-}$, then it follows from (3.8) that we have $\bar{I}_{-} O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+} \in \operatorname{ker} T_{r^{-N}}$. Therefore $\bar{I}_{-} \varphi+\in O_{+} \lambda_{+} \operatorname{ker} T_{r^{-N}}=\operatorname{ker} T_{g}$ and it follows that $\varphi_{+} \in \operatorname{ker} T_{\bar{I}_{-}}$. We conclude that $I_{-}$must be constant since $\operatorname{ker} T_{g}$ is the minimal kernel for $\varphi_{+}$, and $\operatorname{ker} T_{\bar{I}_{-}} \varsubsetneqq \operatorname{ker} T_{g}$ if $I_{-}$is not constant ([7]).

A similar characterisation of non-trivial finite-dimensional Toeplitz kernels in $H^{p}(\mathbb{D})$ was obtained by Nakazi ([20], Theorem 7). The result can be stated as follows (here $\mathbb{T}$ denotes the unit circle):

THEOREM 3.9 ([20]). Let $1<p<\infty$, and $\widetilde{g} \in L_{\infty}(\mathbb{T})$, with $n \in \mathbb{N}$ and let $T_{\widetilde{g}}$ be the associated Toeplitz operator on $H^{p}(\mathbb{D})$. Then the following conditions are equivalent:
(i) $\operatorname{dim} \operatorname{ker} T_{\widetilde{g}}=n$;
(ii) there is a square-rigid outer function $f_{+} \in H^{p}(\mathbb{D})$ such that $\operatorname{ker} T_{\widetilde{g}}=\left\{p f_{+}\right.$: $\left.p \in P_{n-1}\right\}$;
(iii) there is an outer function $h \in H^{\infty}(\mathbb{D})$ with $|\widetilde{g}|=|h|$ on $\mathbb{T}$, and a square-rigid outer function $f_{+} \in H^{p}(\mathbb{D})$ such that

$$
\begin{equation*}
\frac{\tilde{g}}{|\widetilde{g}|} \frac{h}{|h|}=\bar{z}^{n} \frac{\bar{f}_{+}}{f_{+}} \quad \text { on } \mathbb{T} \text {. } \tag{3.9}
\end{equation*}
$$

To compare (3.9) with (3.4), we shall suppose for simplicity that $|g|=1$. Let $m: \mathbb{D} \rightarrow \mathbb{C}_{+}$denote the conformal bijection given by $m(z)=\mathrm{i} \frac{1+z}{1-z}$ for $z \in \mathbb{D}$, and extending this to $\mathbb{T}$, let $\widetilde{g}=g \circ m$. There is a well-known isometric isomorphism $V: L_{p}(\mathbb{T}) \rightarrow L_{p}(\mathbb{R})$, whose restriction maps $H^{p}(\mathbb{D})$ onto $H_{p}^{+}$; it can be defined by

$$
(V \psi)(\xi)=\frac{1}{\pi^{1 / p}} \frac{1}{(\xi+\mathrm{i})^{2 / p}} \psi\left(m^{-1}(\xi)\right), \quad \xi \in \mathbb{R}
$$

However, $V$ does not map the complementary space $\overline{z H^{p}(\mathbb{D})}$ onto $H_{p}^{-}$unless $p=$ 2, so that Toeplitz operators on the disc and half-plane are no longer equivalent. It is therefore not surprising to note that the condition $f_{+} \in H^{p}(\mathbb{D})$ in Nakazi's Theorem 3.9 is not equivalent to the condition $\frac{g_{+}}{\lambda_{+}} \in H_{p}^{+}$in Theorem 3.7 unless $p=2$ 。

Nevertheless, if we define $w(t)=|t-1|^{1-2 / p}$, then we have that the Toeplitz operator $T_{g}$ defined on $H_{p}^{+}$is equivalent to a Toeplitz operator on a weighted

Hardy space $H_{p, w}^{+}$; to define this, let $L_{p, w}(\mathbb{T})=w^{-1} L_{p}(\mathbb{T})$ and note that $B$ : $L_{p}(\mathbb{R}) \rightarrow L_{p, w}(\mathbb{T})$, given by

$$
(B \varphi)(t)=\frac{1}{1-t} \varphi\left(\mathrm{i} \frac{1+t}{1-t}\right), \quad t \in \mathbb{T}
$$

is an isomorphism between $L_{p}(\mathbb{R})$ and $L_{p, w}(\mathbb{T})$ (see [19]). Now let $S_{\mathbb{T}}$ denote the singular integral operator on $L_{p, w}(\mathbb{T})$ defined by

$$
\left(S_{\mathbb{T}} \psi\right)(t)=\frac{1}{\pi \mathrm{i}} \int_{\mathbb{T}} \frac{\psi(\tau)}{\tau-t} \mathrm{~d} \tau, \quad t \in \mathbb{T}
$$

and $H_{p, w}^{ \pm}$the images of the projections $\frac{1}{2}\left(I \pm S_{\mathbb{T}}\right)$. Then $H_{p}^{ \pm}=B^{-1} H_{p, w}^{ \pm}$.
Thus the Toeplitz operator $T_{g}$ on $H_{p}^{+}$is equivalent to a Toeplitz operator $T_{\widetilde{g}}$ on $H_{p, w}^{+}$by $T_{g}=B^{-1} T_{\widetilde{g}} B$.

For $p=2$ we have $w=1$ and there is a unitary equivalence between $T_{g}$ (on $H_{2}^{+}$) and $T_{\widetilde{g}}\left(\right.$ on $H^{2}(\mathbb{D})$ ), so that we recover a case of Theorem 3.9 from our work.

For $p \neq 2$, Theorem 3.7 can be used to extend Nakazi's result to a weighted Hardy space on $\mathbb{D}$; alternatively, Nakazi's result can be used to provide versions of Theorems 3.4 and 3.7 for weighted Hardy spaces of the upper half-plane.

Returning to the half-plane, we have the following.
THEOREM 3.10. If $g \in L_{\infty}$ admits a factorisation

$$
\begin{equation*}
g=g_{-} \theta g_{+}^{-1} \tag{3.10}
\end{equation*}
$$

where $\frac{g_{-}}{\lambda_{-}} \in H_{p}^{-}$is outer, $\frac{g_{+}}{\lambda_{+}} \in H_{p}^{+}$is outer and square-rigid and $\theta \in H_{\infty}^{+}$is an inner function, or if

$$
\begin{equation*}
g=g_{-} r^{N} g_{+}^{-1} \tag{3.11}
\end{equation*}
$$

where $\frac{g_{-}}{\lambda_{-}} \in H_{p}^{-}$is outer and square-rigid, $\frac{g_{+}}{\lambda_{+}} \in H_{p}^{+}$is outer and $N \in \mathbb{N}$, then $\operatorname{ker} T_{g}=\{0\}$.
Proof. Let $O_{+}=\frac{g_{+}^{-1}}{\lambda_{+}}$and $O_{-}=\frac{g_{-}}{\lambda_{-}}$. We have $\varphi_{+} \in \operatorname{ker} T_{g}$ if and only if $\varphi_{+} \in H_{p}^{+}$and $g \varphi_{+}=\varphi_{-}$with $\varphi_{-} \in H_{p}^{-}$which, from 3.10, is equivalent to

$$
O_{-} \lambda_{-} \theta O_{+}^{-1} \lambda_{+}^{-1} \varphi_{+}=\varphi_{-}
$$

Thus if $\varphi_{+} \in \operatorname{ker} T_{g}$ we have

$$
\frac{\bar{O}_{+}}{O_{+}} \frac{\varphi_{+}}{\lambda_{+}}=\frac{\bar{O}_{+} \varphi_{-}}{O_{-}} \lambda_{-}^{-1} \bar{\theta}
$$

where the right hand side represents a function in $H_{p}^{-}$since its conjugate is in $L_{p}$ and in the Smirnov class $\mathcal{N}_{+}$. Therefore $\lambda_{+}^{-1} \varphi_{+} \in \operatorname{ker} T_{\bar{O}_{+} / O_{+}}$and, by near invariance with respect to $\lambda_{+}$([7]), we also have $\varphi_{+} \in \operatorname{ker} T_{\bar{O}_{+} / O_{+}}$. We conclude that $\varphi_{+}=0$ since ker $T_{\bar{O}_{+} / O_{+}}$is one-dimensional.

If $g$ admits a representation (3.11) then $\bar{g}$ admits a factorisation (3.4) and, by Theorem 3.7, $\operatorname{ker} T_{g}^{*}=\operatorname{ker} T_{\bar{g}} \neq\{0\}$. By Coburn's lemma, it follows that $\operatorname{ker} T_{g}=\{0\} . \quad$ ।

It is well known that various properties of a Toeplitz operator can be described in terms of an appropriate factorisation of its symbol. The representations (3.4) and (3.11) generalise the so called $L_{p}$ factorisation, which is a representation of $g$ as a product

$$
\begin{equation*}
g=g_{-} d g_{+}^{-1} \tag{3.12}
\end{equation*}
$$

where

$$
\begin{equation*}
\frac{g_{ \pm}}{\lambda_{ \pm}} \in H_{p}^{ \pm}, \quad \frac{g_{ \pm}^{-1}}{\lambda_{ \pm}} \in H_{q}^{ \pm}, \quad \frac{1}{p}+\frac{1}{q}=1 \tag{3.13}
\end{equation*}
$$

and

$$
\begin{equation*}
d=r^{k}, k \in \mathbb{Z} \tag{3.14}
\end{equation*}
$$

The Toeplitz operator $T_{g}$ is Fredholm if and only if $g$ admits a factorisation (3.12) satisfying the conditions (3.13), (3.14) and such that $g_{-} P^{+} g_{-}^{-1} I$ is a densely defined bounded operator in $L_{p}$. In that case 3.12 is called a generalised $p$-factorisation, or Wiener-Hopf factorisation relative to $L_{p}$, and we have $\operatorname{dim} \operatorname{ker} T_{g}=k$ if $k \geqslant 0$, dim ker $T_{g}^{*}=-k$ if $k \leqslant 0$ ([5], [18], [19]).

As an illustration, we consider $g(\xi)=r^{\alpha}, \alpha \in\left[-\frac{1}{2}, \frac{1}{2}\right]$, where we assume the discontinuity to be at $\infty$. A Wiener-Hopf factorisation relative to $L_{2}$ exists for all $\alpha \neq \pm \frac{1}{2}$ and $T_{g}$ is invertible in $H_{2}^{+}$([11]). If $\alpha= \pm \frac{1}{2}$ then $g$ does not admit a Wiener-Hopf factorisation relative to $L_{2}$. However we can write $g=g_{-} r^{N} g_{+}^{-1}$ with $g_{-}=(\xi-i)^{-1 / 2}, g_{+}^{-1}=(\xi+i)^{1 / 2}$ and $N=0$ if $\alpha=-\frac{1}{2}, N=1$ if $\alpha=\frac{1}{2}$. It is clear that $g_{ \pm} \lambda_{ \pm}^{-1} \in H_{2}^{ \pm}$are outer and, since we have

$$
\mathcal{K}_{\min }\left(g_{+} \lambda_{+}^{-1}\right)=\operatorname{ker} T_{r^{-3 / 2}}=\operatorname{span}\left\{g_{+} \lambda_{+}^{-1}\right\}
$$

the function $g_{+} \lambda_{+}^{-1}$ is square-rigid. We have thus $\operatorname{ker} T_{g}=\operatorname{ker} T_{g}^{*}=\{0\}$, in accordance with Theorem 3.10.

The representation (3.12) is called a bounded factorisation if $g_{-}^{ \pm 1} \in H_{\infty}^{-}$and $g_{+}^{ \pm 1} \in H_{\infty}^{+}([4])$. In various subalgebras of $L_{\infty}$, every invertible element admits a bounded factorisation (3.12 where $d$ is an inner function. This is the case for the Wiener algebra and the algebra of all Hölder continuous functions with exponent $\mu \in(0,1)$, with $d=r^{k}, k \in \mathbb{Z}([18],[19])$, and the algebra $A P$ of almost periodic functions, with $d(\xi)=\exp (-i \lambda \xi), \lambda \in \mathbb{R}([10],[14])$.

In the latter case, we easily see moreover that, for every $g \in A P$ which is invertible in $L_{\infty}$ (and thus also in $A P$ ), $\operatorname{ker} T_{g}$ is either trivial or isomorphic to an infinite dimensional model space $K_{\theta}$ with $\theta(\xi)=\exp (\mathrm{i} \lambda \xi)$, depending on whether $\lambda \leqslant 0$ or $\lambda \geqslant 0$.

Acknowledgements. This research was partially supported by FCT/Portugal through Projects PTDC/MAT/121837/2010 and UID/MAT/04459/2013.

## REFERENCES

[1] S. Barclay, A solution to the Douglas-Rudin problem for matrix-valued functions, Proc. London Math. Soc. (3) 99(2009), 757-786.
[2] C. Benhida, M.C. CÂmara, C. Diogo, Some properties of the kernel and the cokernel of Toeplitz operators with matrix symbols, Linear Algebra Appl. 432(2010), 307-317.
[3] J. Bourgain, A problem of Douglas and Rudin on factorization, Pacific J. Math. 121(1986), 47-50.
[4] M.C. Câmara, C. Diogo, Yu.I. Karlovich, I.M. SpitKovsky, Factorizations, Riemann-Hilbert problems and the corona theorem, J. London Math. Soc. (2) 86(2012), 852-878.
[5] M.C. Câmara, C. Diogo, L. Rodman, Fredholmness of Toeplitz operators and corona problems, J. Funct. Anal. 259(2010), 1273-1299.
[6] M.C. Câmara, M.T. Malheiro, J.R. Partington, Model spaces in reflexive Hardy spaces, Operators and Matrices, to appear.
[7] M.C. CÂmara, J.R. Partington, Near invariance and kernels of Toeplitz operators, J. Anal. Math. 124(2014), 235-260.
[8] I. Chalendar, N. Chevrot, J.R. Partington, Nearly invariant subspaces for backwards shifts on vector-valued Hardy spaces, J. Operator Theory 63(2010), 403-415.
[9] N. Chevrot, Kernel of vector-valued Toeplitz operators, Integral Equations Operator Theory 67(2010), 57-78.
[10] L. Coburn, R.G. Douglas, Translation operators on the half-line, Proc. Nat. Acad. Sci. USA 62(1969), 1010-1013.
[11] R. Duduchava, Integral Equations in Convolution with Discontinuous Presymbols, Singular Integral Equations with Fixed Singularities, and their Applications to some Problems of Mechanics, Teubner-Texte zur Mathematik, BSB B. G. Teubner Verlagsgesellschaft, Leipzig 1979.
[12] P.L. Duren, Theory of $H^{p}$ Spaces, Dover, New York 2000.
[13] K.M. DYaKonov, Kernels of Toeplitz operators via Bourgain's factorization theorem, J. Funct. Anal. 170(2000), 93-106.
[14] I.C. Gohberg, I.A. Feldman, Wiener-Hopf integro-difference equations [Russian], Dokl. Akad. Nauk SSSR 183(1968), 25-28; English translation: Soviet Math. Dokl. 9(1968), 1312-1316.
[15] E. Hayashi, The kernel of a Toeplitz operator, Integral Equations Operator Theory 9(1986), 588-591.
[16] E. HAYASHI, Classification of nearly invariant subspaces of the backward shift, Proc. Amer. Math. Soc. 110(1990), 441-448.
[17] D. Hitt, Invariant subspaces of $\mathcal{H}^{2}$ of an annulus, Pacific J. Math. 134(1988), 101-120.
[18] G.S. Litvinchuk, I.M. Spitkovskii, Factorization of Measurable Matrix Functions, Oper. Theory Adv. Appl., vol. 25, Birkhäuser Verlag, Basel 1987.
[19] S.G. Mikhlin, S. PrösSdorf, Singular Integral Operators, Springer-Verlag, Berlin 1986.
[20] T. NAKAZI, Kernels of Toeplitz operators, J. Math. Soc. Japan 38(1986), 607-616.
[21] D. Sarason, Nearly invariant subspaces of the backward shift, in Contributions to Operator Theory and its Applications (Mesa, AZ, 1987), Oper. Theory Adv. Appl., vol. 35, Birkhäuser, Basel 1988, pp. 481-493.
[22] D. Sarason, Kernels of Toeplitz operators, in Toeplitz Operators and Related Topics (Santa Cruz, CA, 1992), Oper. Theory Adv. Appl., vol. 71, Birkhäuser, Basel 1994, pp. 153-164.
[23] H. Widom, Singular integral equations in $L_{p}$, Trans. Amer. Math. Soc. 97(1960), 131160.
M.C. CÂMARA, CEnter for Mathematical Analysis, Geometry and Dynamical Systems, Mathematics Department, Instituto Superior Técnico, UniVersidade de Lisboa, Av. Rovisco Pais, 1049-001 Lisboa, Portugal

E-mail address: ccamara@math.tecnico.ulisboa.pt
J.R. PARTINGTON, SChool of Mathematics, University of Leeds, Leeds LS2 9JT, U.K.

E-mail address: j.r.partington@leeds.ac.uk

